### NEW INTEGRABILITY CONDITIONS OF DERIVATIONAL EQUATIONS OF A SUBMANIFOLD IN A GENERALIZED RIEMANNIAN SPACE

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#### Abstract

The present work is a continuation of [5] and [6]. In [5] we have obtained derivational equations of a submanifold  $X_M$  of a generalized Riemannian space  $GR_N$ . Since the basic tensor in  $GR_N$  is asymmetric and in this way the connection is also asymmetric, in a submanifold the connection is generally asymmetric too. By reason of this, we define 4 kinds of covariant derivative and obtain 4 kinds of derivational equations. In [6] we have obtained integrability conditions and Gauss-Codazzi equations using the  $1^{st}$  and the  $2^{st}$  kind of covariant derivative.

The present work deals in the cited matter, using the  $3^{rd}$  and the  $4^{th}$  kind of covariant derivative. One obtains three new integrability conditions for derivational equations of tangents and three such conditions for normals of the submanifold, as the corresponding Gauss-Codazzi equations too.

### 1 Introduction

**1.1.** A generalized Riemannian space  $GR_N$  is a differentiable manifold equipped with an asymmetric basic tensor  $G_{ij}(x^1,...,x^N)$  (the components) where  $x^i$  are the local coordinates. The symmetric, respectively antisymmetric part of  $G_{ij}$  are  $H_{ij}$  and  $K_{ij}$ .

For the lowering and rasing of indices in  $GR_N$  one uses  $H_{ij}$ , respectively  $H^{ij}$ , where

(1.1) 
$$(H^{ij}) = (H_{ij})^{-1}, \quad (det(H_{ij}) \neq 0).$$

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Cristoffel symbols at  $GR_N$  are

(1.2) 
$$\Gamma_{i,jk} = \frac{1}{2} (G_{ji,k} - G_{jk,i} + G_{ik,j}), \quad \Gamma_{jk}^i = H^{ip} \Gamma_{p,jk},$$

where, for example,  $G_{ji,k} = \partial G_{ji}/\partial x^k$ . Based on the asymmetry of  $G_{ij}$ , it follows that the Cristoffel symbols are also asymmetric with respect to j, k in (1.2).

By equations

(1.3) 
$$x^{i} = x^{i}(u^{1}, ..., u^{M}) \equiv x^{i}(u^{\alpha}), \quad i = 1, ..., N,$$

a submanifold  $X_M$  is defined in local coordinates. If  $rank(B^i_\alpha)=M$   $(B^i_\alpha=\partial x^i/\partial u^\alpha)$  and

$$(1.4) g_{\alpha\beta} = B_{\alpha}^{i} B_{\beta}^{j} G_{ij},$$

 $X_M$  becomes  $GR_M \subset GR_N$ , with **induced basic tensor** (1.4), which is generally also asymmetric. Note that in the present work Latin indices i, j, ... take values 1, ..., N and refer to the  $GR_N$ , while the Greek ones take values 1, ..., M and refer to the  $GR_M$ .

In the  $GR_M$  are valid the relations similar to (1.1) and (1.2). The symmetric part of  $g_{\alpha\beta}$  is denoted with  $h_{\alpha\beta}$ , and antisymmetric one with  $k_{\alpha\beta}$ , where e.g.

$$(1.5) h_{\alpha\beta} = B_{\alpha}^i B_{\beta}^j H_{ij}, \quad (h^{\alpha\beta}) = (h_{\alpha\beta})^{-1}.$$

Cristoffel symbols  $\widetilde{\Gamma}_{\alpha.\beta\gamma}$ ,  $\widetilde{\Gamma}^{\alpha}_{\beta\gamma} = h^{\alpha\pi}\widetilde{\Gamma}_{\pi.\beta\gamma}$  are expressed by  $g_{\alpha\beta}$  analogously to (1.2).

For the unit, mutually orthogonal vectors  $N_A^i$ , which are orthogonal to the  $GR_M$  too, we have [1]

(1.6) 
$$H_{ij}N_A^iN_B^j = e_A\delta_B^A = h_{AB}, \ e_A \in \{-1,1\}, \ H_{ij}N_A^iB_\alpha^j = 0,$$

where  $A, B, \dots \in \{M+1, \dots, N\}$ .

As it is known, the following relations between Cristoffel symbols of a generalized Riemannian space and its subspace are valid:

(1.7) 
$$\widetilde{\Gamma}_{\alpha,\beta\gamma} = \Gamma_{i,jk} B^i_{\alpha} B^j_{\beta} B^k_{\gamma} + H_{ij} B^i_{\alpha} B^j_{\beta,\gamma},$$

(1.8) 
$$\widetilde{\Gamma}^{\alpha}_{\beta\gamma} = h^{\pi\alpha}\widetilde{\Gamma}_{\pi.\beta\gamma} = h^{\pi\alpha}(\Gamma_{i.jk}B^i_{\pi}B^j_{\beta}B^k_{\gamma} + H_{ij}B^i_{\pi}B^j_{\beta,\gamma}),$$

i.e.

$$\widetilde{\Gamma}^{\alpha}_{\beta\gamma} = h^{\pi\alpha} H_{pi} B^p_{\pi} (\Gamma^i_{jk} B^j_{\beta} B^k_{\gamma} + B^i_{\beta,\gamma}).$$

**1.2.** The set of normals of the submanifold  $X_M \subset GR_N$  make a **normal bundle** for  $X_M$ , and we note it  $X_{N-M}^N$ . One can introduce a metric tensor on  $X_{N-M}^N$ 

$$g_{AB} = G_{ij} N_A^i N_B^j,$$

which is asymmetric in a general case.

The symmetric part is

(1.10) 
$$h_{AB} = H_{ij} N_A^i N_B^j = e_A \delta_B^A = h_{BA} = \begin{cases} e_A, & A=B, \\ 0, & \text{otherwise.} \end{cases}, e_A \in \{-1, 1\}.$$

If

$$(h^{AB}) = (h_{AB})^{-1},$$

we have

$$h^{AB} = e_A \delta_B^A = h_{AB} = h^{BA}.$$

On  $X_{N-M}^N$  one can define in two manners connection coefficients

(1.11) 
$$\overline{\Gamma}_{B\mu}^{A} = H_{ij} h^{AQ} N_Q^j (N_{B,\mu}^i + \Gamma_{pq}^i N_B^p B_\mu^q).$$

Being the coefficients  $\Gamma$ ,  $\widetilde{\Gamma}$ ,  $\overline{\Gamma}$  non-symmetric in general, for a tensor, defined at points of  $GR_M$ , is possible define four kinds of covariant derivative. For example

$$(1.12) \begin{array}{c} \nabla_{\mu}t_{j\beta B}^{i\alpha A}\equiv t_{j\beta B}^{i\alpha A}|_{\mu}=t_{j\beta B,\mu}^{i\alpha A}+\Gamma_{pm}^{i}t_{j\beta B}^{p\alpha A}B_{\mu}^{m}-\Gamma_{jm}^{p}t_{p\beta B}^{i\alpha A}B_{\mu}^{m}\\ \frac{1}{2}\\ \frac{1}{3}\\ \frac{2}{3}\\ \frac{2}{3}\\ \frac{2}{3}\\ \frac{2}{3}\\ \frac{2}{3}\\ \frac{2}{3}\\ \frac{2}{mp}\\ \frac{pm}{mp}\\ \frac{pm}{mp}\\ \frac{mj}{jm}\\ \frac{mj}{jm}\\ \frac{mj}{jm}\\ \frac{mj}{jm}\\ \frac{mj}{jm}\\ \frac{pm}{mp}\\ \frac{p$$

In this way four connection  $\nabla_{\theta}$ ,  $\theta \in \{1, \dots, 4\}$ , on  $X_M \subset GR_N$  are defined. We shall note the obtained structures  $(X_M \subset GR_N, \nabla_{\theta}, \theta \in \{1, \dots, 4\})$ .

## 2 New first and second kind integrability conditions of derivational equations

**2.0.** In [5] are obtained derivational equations of a submanifold in a  $GR_N$ , and in [6] integrability conditions of these equations in the structure  $(X_M \subset GR_N, \nabla_{\theta}, \theta \in \{1,2\})$ . In the present work we engage in this problem for the structure  $(X_M \subset GR_N, \nabla_{\theta}, \theta \in \{3,4\})$ .

As it is proved in [5] (Th. 1.2.), derivational equations in the considered case for a tangent are

(2.1) 
$$B_{\alpha|\mu}^{i} = \sum_{P} \Omega_{P\alpha\mu} N_{P}^{i}, \quad \theta \in \{3, 4\},$$

and then for induced torsion in  $X_M$  is valid

(2.2) 
$$\widetilde{T}^{\alpha}_{\beta\gamma} = 0 \ (\widetilde{\Gamma}^{\alpha}_{\beta\gamma} = \widetilde{\Gamma}^{\alpha}_{\gamma\beta}).$$

By virtue of the Th. 2.3. in [5], for unit normal is

$$(2.3) N_{A \mid \mu}^i = -e_A \Omega_{A\rho\mu} h^{\pi\rho} B_\pi^i, \quad \theta \in \{3,4\},$$

and

(2.4) 
$$\overline{\Gamma}_{1B\mu}^{A} = \overline{\Gamma}_{2B\mu}^{A} = \overline{\Gamma}_{B\mu}^{A},$$

in (1.12), and based on (1.8) in [5]

(2.5) 
$$\Omega_{P\alpha\mu} = e_P H_{ij} N_P^i (B_{\alpha,\mu}^j + \Gamma_{pm}^j B_{\alpha}^p B_{\mu}^m) = \Omega_{P\alpha\mu}.$$

In relation with (2.2,4), the addends in (1.12), related to  $X_M$  and to  $X_{N-M}^N$  are not different for separate kinds of derivatives, and (1.12) now becomes

$$(2.6) \begin{array}{c} t_{j\beta B|\mu}^{i\alpha A} = t_{j\beta B,\mu}^{i\alpha A} + \Gamma_{pm}^{i} t_{j\beta B}^{p\alpha A} B_{\mu}^{m} - \Gamma_{jm}^{p} t_{p\beta B}^{i\alpha A} B_{\mu}^{m} \\ \frac{2}{3} & \frac{pm}{mp} & \frac{mj}{jm} \\ + \widetilde{\Gamma}_{\pi\mu}^{\alpha} t_{j\beta B}^{i\pi A} - \widetilde{\Gamma}_{\beta\mu}^{\pi} t_{j\pi B}^{i\alpha A} + \overline{\Gamma}_{P\mu}^{A} t_{j\beta B}^{i\alpha P} - \overline{\Gamma}_{B\mu}^{P} t_{j\beta P}^{i\alpha A}, \end{array}$$

where the coefficients  $\widetilde{\Gamma}$  are symmetric, and  $\overline{\Gamma}$  are unique  $(\overline{\Gamma}_1 = \overline{\Gamma}_2 = \overline{\Gamma})$ . If in a differentiated tensor no exists indices as i, j, ..., we write  $|\mu$  instead of  $\mu$ .

Using (2.1,3), we get (see (2.4) in [6])

(2.7) 
$$B_{\alpha|\mu|\nu}^{i} - B_{\alpha|\nu|\mu}^{i} = \sum_{P} [e_{P}h^{\pi\rho}(-\Omega_{P\alpha\mu}\Omega_{\omega}P_{\rho\nu} + \Omega_{P\alpha\nu}\Omega_{\theta}P_{\rho\mu})B_{\pi}^{i} + (\Omega_{P\alpha\mu|\nu} - \Omega_{P\alpha\mu|\mu}N_{\rho}^{i})N_{\rho}^{i}], \quad \theta, \omega \in \{3, 4\}.$$

**2.1.** With respect of Ricci-type identities (12) and (13) from [2], and taking into consideration (2.2), we have

$$(2.8) B_{\alpha \mid \mu \mid \nu}^{i} - B_{\alpha \mid \nu \mid \mu}^{i} = \underset{\theta \rightarrow 2}{R^{i}_{pmn}} B_{\alpha}^{p} B_{\mu}^{m} B_{\nu}^{n} - \widetilde{R}_{\alpha \mu \nu}^{\pi} B_{\pi}^{i}, \ \theta \in \{3, 4\},$$

where

$$(2.9a) R_{jmn}^i = \Gamma_{jm,n}^i - \Gamma_{jn,m}^i + \Gamma_{jm}^p \Gamma_{pn}^i - \Gamma_{jn}^p \Gamma_{pm}^i,$$

$$(2.9b) R_{jmn}^i = \Gamma_{mj,n}^i - \Gamma_{nj,m}^i + \Gamma_{mj}^p \Gamma_{np}^i - \Gamma_{nj}^p \Gamma_{mp}^i$$

are curvature tensors of the 1<sup>st</sup>, respectively 2<sup>nd</sup> kind of  $GR_N$  and  $\widetilde{R}^{\alpha}_{\beta\mu\nu}$  is, with respect of (2.2), curvature tensor of  $R_M \subset GR_N$ .

We obtained in [6] three kinds integrability conditions for derivational equation of a tangent  $B^i_{\alpha}$ , i.e. for  $B^i_{\alpha|\mu}$ ,  $\theta \in \{1,2\}$ . We shall consider here such conditions for  $\theta \in \{3,4\}$ .

If one substitutes  $\theta = \omega \in \{3, 4\}$  into (2.7) and compares with (2.8), taking into consideration (2.5) and (2.6), we get

(2.10) 
$$R_{\alpha}^{i} B_{\alpha}^{p} B_{\mu}^{m} B_{\nu}^{n} = \left[ \widetilde{R}_{\alpha\mu\nu}^{\pi} - \sum_{P} e_{P} h^{\pi\rho} \left( \Omega_{P\alpha\mu} \Omega_{P\rho\nu} - \Omega_{P\alpha\nu} \Omega_{P\rho\mu} \right) \right] B_{\pi}^{i} + \sum_{P} \left[ \Omega_{P\alpha\mu|\nu} - \Omega_{P\alpha\nu|\mu} \right] N_{P}^{i}, \quad \theta \in \{3, 4\},$$

which are the 1<sup>st</sup> and the 2<sup>nd</sup> integrability conditions of derivational equation (2.1) in the structure  $(X_M \subset GR_N, \ \nabla_{\alpha}, \ \theta \in \{3,4\})$ .

a) Composing the previous equation with  $H^{ij}B^j_{\beta}$ , one gets

$$(2.11) \ \ \underset{\theta = 2}{R} jpmn B^{j} \beta B^{p}_{\alpha} B^{m}_{\mu} B^{n}_{\nu} = \widetilde{R}_{\beta\alpha\mu\nu} - \sum_{P} e_{P} (\underset{\theta}{\Omega_{P\alpha\mu}} \underset{\theta}{\Omega_{P\beta\nu}} - \underset{\theta}{\Omega_{P\alpha\nu}} \underset{\theta}{\Omega_{P\beta\mu}}), \ \theta \in \{3,4\},$$

where

$$(2.12 a, b) R_{\theta \to jpmn} = H_{ij} R_{\theta \to pmn}^{i}, \widetilde{R}_{\beta\alpha\mu\nu} = h_{\pi\beta} \widetilde{R}_{\alpha\mu\nu}^{\pi}, \theta \in \{3, 4\}.$$

Taking into count the antisymmetry of the tensors (2.12) with respect of the first two indices and substituting i in place of p, the equation (2.11) becomes

$$(2.13) \quad \widetilde{R}_{\alpha\beta\mu\nu} = \underset{\theta \rightarrow 2}{R}_{ijmn} B^i_{\alpha} B^j_{\beta} B^m_{\mu} B^n_{\nu} - \sum_{P} e_P (\underset{\theta}{\Omega_{P\alpha\mu}} \underset{\theta}{\Omega_{P\beta\nu}} - \underset{\theta}{\Omega_{P\alpha\nu}} \underset{\theta}{\Omega_{P\beta\mu}}), \ \theta \in \{3,4\},$$

which are Gauss equations of the 1<sup>st</sup> and the 2<sup>nd</sup> kind in the structure  $(X_M \subset GR_N, \nabla_{\theta}, \theta \in \{3,4\})$ .

b) Composing the equation (2.10) with  $H_{ij}N_Q^j$  we obtain finally

(2.14) 
$$R_{\alpha P_{ijmn}}B_{\alpha}^{i}N_{Q}^{j}B_{\mu}^{m}B_{\nu}^{n} = e_{Q}(\Omega_{Q\alpha\nu|\mu} - \Omega_{Q\alpha\mu|\nu}), \ \theta \in \{3,4\},$$

and that are the  $1^{st}$  Codazzi equations of the  $1^{st}$  and the  $2^{nd}$  kind at the cited structure.

**2.2.** Consider the same matter for the unit normal  $N_A^i$ . Using (2.3,1), we obtain (see (2.13) in [6]):

$$(2.15) N_{A_{\omega}^{i} \mu_{\omega}^{i} \nu}^{i} - N_{A_{\omega}^{i} \nu_{\theta}^{i} \mu}^{i} = -e_{A} h^{\pi \rho} [(\Omega_{A \rho \mu_{\omega}^{i} \nu} - \Omega_{\omega}^{i} A_{\rho \nu_{\theta}^{i} \mu}) B_{\pi}^{i} + \sum_{P} (\Omega_{A \rho \mu}^{i} \Omega_{P \pi \nu} - \Omega_{\omega}^{i} A_{\rho \nu} \Omega_{\theta}^{i} P_{\pi \mu}) N_{P}^{i}].$$

In order to find corresponding Ricci-type identity for the left side of this equation for  $\theta = \omega \in \{3, 4\}$ , we use (2.6). Firstly, we have

(2.16) 
$$N_{A|\mu}^{i} = N_{A,\mu}^{i} + \Gamma_{pm}^{i} N_{A}^{p} B_{\mu}^{m} - \overline{\Gamma}_{A\mu}^{p} N_{P}^{i},$$

and further

$$\begin{split} N_{A|\mu|\nu}^i &= (N_{A|\mu}^i)_{,\nu} + \Gamma_{sn}^i N_{A|\mu}^s B_{\nu}^n - \widetilde{\Gamma}_{\mu\nu}^{\sigma} N_{A|\sigma}^s - \overline{\Gamma}_{A\nu}^S N_{S|\mu}^i \\ &= N_{A,\mu\nu}^i + \Gamma_{pm,n}^i N_A^p B_{\mu}^m B_{\nu}^n + \Gamma_{pm}^i N_{A,\nu}^p B_{\mu}^m + \Gamma_{pm}^i N_A^p B_{\mu,\nu}^m \\ &- \overline{\Gamma}_{A\mu,\nu}^P N_P^i - \overline{\Gamma}_{A\mu}^P N_{P,\nu}^i + \Gamma_{sn}^i N_{A,\mu}^s B_{\nu}^n + \Gamma_{sn}^i \Gamma_{pm}^s B_{\nu}^n N_A^p B_{\mu}^m \\ &- \Gamma_{sn}^i N_P^s \overline{\Gamma}_{A\mu}^P B_{\nu}^n - \widetilde{\Gamma}_{\mu\nu}^{\sigma} N_{A,\sigma}^i - \widetilde{\Gamma}_{\mu\nu}^{\sigma} \Gamma_{pm}^i N_A^p B_{\sigma}^m + \widetilde{\Gamma}_{\mu\nu}^{\sigma} \overline{\Gamma}_{A\sigma}^P N_P^i \\ &- \overline{\Gamma}_{A\nu}^S N_{S,\mu}^i - \overline{\Gamma}_{A\nu}^S \Gamma_{pm}^i N_S^p B_{\mu}^m + \overline{\Gamma}_{A\nu}^S \overline{\Gamma}_{S\mu}^P N_P^i, \end{split}$$

wherefrom

$$(2.17) N_{A_{\alpha}\mu}^{i}{}_{\alpha}{}^{\nu} - N_{A_{\alpha}\nu}^{i}{}_{\alpha}{}^{\mu} = \overline{R}_{pmn}^{i} N_{A}^{p} B_{\mu}^{m} B_{\nu}^{n} - \overline{R}_{A\mu\nu}^{P} N_{P}^{i}.$$

where

$$\overline{R}_{B\mu\nu}^{A} = \overline{\Gamma}_{B\mu,\nu}^{A} - \overline{\Gamma}_{B\nu,\mu}^{A} + \overline{\Gamma}_{B\mu}^{P} \overline{\Gamma}_{P\nu}^{A} - \overline{\Gamma}_{B\nu}^{P} \overline{\Gamma}_{P\mu}^{A},$$

is curvature tensor of the spaceGR<sub>N</sub> with respect to the normal submanifold in the structure  $(X_M \subset GR_N, \ \nabla_{\alpha}, \ \theta \in \{3,4\})$ .

By means of the  $4^{th}$  kind of covariant derivative we obtain an equation corresponding to (2.17), and we conclude

$$(2.19) N_{A|\mu|\nu}^{i} - N_{A|\nu|\mu}^{i} = \underset{\theta-2}{R}_{pmn}^{i} N_{A}^{p} B_{\mu}^{m} B_{\nu}^{n} - \overline{R}_{A\mu\nu}^{P} N_{P}^{i}, \ \theta \in \{3, 4\}.$$

If one substitutes into (2.15)  $\theta = \omega \in \{3,4\}$  and equilizes the right sides of obtained equation and (2.19), we get **the 1<sup>st</sup> and the 2<sup>nd</sup> kind integrability conditions of derivational equation (2.3)** in the structure  $(X_M \subset GR_N, \nabla_{\theta}, \theta \in \{3,4\})$ :

$$(2.20) \qquad \begin{aligned} R_{\theta-2}^{i} N_{A}^{p} B_{\mu}^{m} B_{\nu}^{n} &= e_{A} h^{\pi \rho} (\Omega_{\theta}^{i} A_{\rho \mu} |_{\nu} - \Omega_{\theta}^{i} A_{\rho \nu} |_{\mu}) B_{\pi}^{i} \\ &+ [\overline{R}_{A\mu\nu}^{P} - e_{A} h^{\pi \rho} \sum_{P} (\Omega_{\theta}^{i} A_{\rho \mu} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} \Omega_{P}^{i} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} \Omega_{P}^{i} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} \Omega_{P}^{i} \Omega_{P}^{i} \Omega_{P}^{i} + \Omega_{\theta}^{i} \Omega_{P}^{i} \Omega_{P}^{i$$

- a) If we compose this equation with  $H_{ij}B^j_{\beta}$  one obtains an equation equivalent with (2.14),that is the 1<sup>st</sup> Codazzi equation of the 1<sup>st</sup> and the 2<sup>nd</sup> kind for the structure  $(X_M \subset GR_N, \ \nabla_{\alpha}, \ \theta \in \{3,4\}).$ 
  - b) By composing the equation (2.20) with  $H_{ij}N_B^j$ , one obtains endly

$$(2.21) \qquad \underset{\theta \to 2}{R_{ijmn}} N_A^i N_B^j B_\mu^m B_\nu^n = \overline{R}_{AB\mu\nu} + e_A e_B h^{\pi\rho} (\underset{\theta}{\Omega}_{A\pi\mu} \underset{\theta}{\Omega}_{B\rho\nu} - \underset{\theta}{\Omega}_{A\pi\nu} \underset{\theta}{\Omega}_{B\rho\mu}),$$

where

$$\overline{R}_{AB\mu\nu} = h_{AP} \overline{R}_{B\mu\nu}^P.$$

The equation (2.21) is the **2**<sup>nd</sup> Codazzi equation of the **1**<sup>st</sup> and the **2**<sup>nd</sup> kind for the structure  $(X_M \subset GR_N, \nabla_{\alpha}, \theta \in \{3,4\})$ .

Based on expressed above, the next theorems are valid:

**Theorem 2.1.** The 1<sup>st</sup> and the 2<sup>nd</sup> kind integrability conditions for derivational equations (2.1), (2.3) in the in the structure  $(X_M \subset GR_N, \nabla_{\theta}, \theta \in \{3,4\})$  are given by equations (2.10), (2.20) respectively, where  $\Omega$  is given in (2.5), R, R in (2.9), R is curvature tensor of the symmetric connection  $\Gamma$ , while R is given in (2.18), (2.22).

**Theorem 2.2.** The Gauss equations of the 1<sup>st</sup> and the 2<sup>nd</sup> kind in the structure  $(X_M \subset GR_N, \nabla, \theta \in \{3,4\})$  are given in (2.13), the 1<sup>st</sup> Codazzi equations of the 1<sup>st</sup> and the 2<sup>nd</sup> kind in (2.14), and the 2<sup>nd</sup> Codazzi equations of the 1<sup>st</sup> and the 2<sup>nd</sup> kind in (2.21) in the same structure.

# 3 Third kind integrability condition of derivational equations

**3.1.** Using simultaneously the  $3^{rd}$  and the  $4^{th}$  kind of covariant derivative by virtue of (2.6), we obtain Ricci-type identity (eq. (46) in [2]):

$$(3.1) B_{\alpha_{|\mu|\nu}^{i}}^{i} - B_{\alpha_{|\nu|\mu}^{i}}^{i} = R_{4p\mu\nu}^{i} B_{\alpha}^{p} - \widetilde{R}_{\alpha\mu\nu}^{\pi} B_{\pi}^{i},$$

where

$$(3.2) \qquad R^{i}_{_{A}j\mu\nu} = (\Gamma^{i}_{jm,n} - \Gamma^{i}_{nj,m} + \Gamma^{p}_{jm}\Gamma^{i}_{np} - \Gamma^{p}_{nj}\Gamma^{i}_{pm})B^{m}_{\mu}B^{n}_{\nu} + T^{i}_{jm}(B^{m}_{\mu,\nu} - \widetilde{\Gamma}^{\pi}_{\nu\mu}B^{m}_{\pi})$$

is curvature tensor of the  $4^{th}$  kind of  $GR_N$  with respect to  $X_M\subset GR_N.$ 

On the other hand, if we put into (2.7)  $\theta = 3$ ,  $\omega = 4$  and compare the obtained equation with (3.1), we obtain **the 3<sup>rd</sup> kind integrability condition** of derivational equation (2.1) in the structure  $(X_M \subset GR_N, \nabla, \theta \in \{3,4\})$ :

a) Composing previous equation with  $H_{ij}B^j_{\beta}$ , we get

$$R_{4jp\mu\nu}B_{\beta}^{j}B_{\alpha}^{p} = \widetilde{R}_{\beta\alpha\mu\nu} - \sum_{P} e_{P} \left( \Omega_{P\alpha\mu}\Omega_{P\beta\nu} - \Omega_{P\alpha\nu}\Omega_{P\beta\mu} \right),$$

i.e., exchanging  $j \to i, \, p \to j, \, \alpha \leftrightarrow \beta$ , it follows that

$$\widetilde{R}_{\alpha\beta\mu\nu} = R_{ij\mu\nu} B_{\alpha}^{i} B_{\beta}^{j} - \sum_{P} e_{P} \left( \Omega_{P\alpha\mu} \Omega_{P\beta\nu} - \Omega_{P\alpha\nu} \Omega_{P\beta\mu} \right),$$

where

$$R_{ij\mu\nu} = H_{ip}R_{j\mu\nu}^p.$$

The equation (3.4) is Gauss equation of the 3<sup>rd</sup> in the structure  $(X_M \subset GR_N, \nabla_{\underline{a}}, \theta \in \{3,4\}).$ 

b) Composing (3.4) with  $H_{ij}N_Q^j$ , we obtain

$$R_{ij\mu\nu}N_Q^iB_\alpha^j = e_Q(\Omega_{Q\alpha\mu|\nu} - \Omega_{Q\alpha\nu|\mu}).$$

This is the 1<sup>st</sup> Codazzi equation of the 3<sup>rd</sup> kind in the cited structure.

**3.2.** On the base of (2.6) and (2.16) we have

$$\begin{split} N_{A|\mu|\nu}^i &= (N_{A|\mu}^i)_{,\nu} + \Gamma_{ns}^i N_{A|\mu}^s B_{\nu}^n - \widetilde{\Gamma}_{\mu\nu}^{\sigma} N_{A|\sigma}^i - \overline{\Gamma}_{A\nu}^S N_{S|\mu}^i \\ &= N_{A,\mu\nu}^i + \Gamma_{pm,n}^i N_A^p B_{\mu}^m B_{\nu}^n + \Gamma_{pm}^i N_{A,\nu}^p B_{\mu}^m + \Gamma_{pm}^i N_A^p B_{\mu,nu}^m \\ &- \overline{\Gamma}_{A\mu,\nu}^P N_P^i - \overline{\Gamma}_{A\mu}^P N_{P,\nu}^i + \Gamma_{ns}^i N_{A,\mu}^s B_{\nu}^n + \Gamma_{ns}^i \Gamma_{pm}^s B_{\nu}^n N_A^p B_{\mu}^m \\ &- \Gamma_{ns}^i N_P^s \overline{\Gamma}_{A\mu}^P B_{\nu}^n - \widetilde{\Gamma}_{\mu\nu}^\sigma N_{A,\sigma}^i - \widetilde{\Gamma}_{\mu\nu}^\sigma \Gamma_{pm}^i N_A^p B_{\sigma}^m + \widetilde{\Gamma}_{\mu\nu}^\sigma \overline{\Gamma}_{A\sigma}^P N_P^i \\ &- \overline{\Gamma}_{A\nu}^S N_{S,\mu}^i - \overline{\Gamma}_{A\nu}^S \Gamma_{mm}^i N_P^s B_{\mu}^m + \overline{\Gamma}_{A\nu}^S \overline{\Gamma}_{S\mu}^P N_P^i, \end{split}$$

and

$$(3.6) N_{A \mid \mu \mid \nu}^{i} - N_{A \mid \nu \mid \mu}^{i} = R_{4p\mu\nu}^{i} N_{A}^{p} - \overline{R}_{A\mu\nu}^{P} N_{P}^{i},$$

where R is given in (3.2), and  $\overline{R}$  in (2.18).

By substituting into (2.15)  $\theta = 3$ ,  $\omega = 4$  and comparing the obtained equation with (3.6), we obtain the  $3^{rd}$  kind integrability condition of derivational **equation (2.3)** in the structure  $(X_M \subset GR_N, \nabla, \theta \in \{3,4\})$ :

(3.11) 
$$R_{4p\mu\nu}^{i}N_{A}^{p} = -e_{A}h^{\pi\rho}(\Omega_{A\rho\mu|\nu} - \Omega_{A\rho\nu|\mu})B_{\pi}^{i} + [\overline{R}_{A\mu\nu}^{P} - e_{A}h^{\pi\rho}\sum_{P}(\Omega_{A\rho\mu}\Omega_{P\pi\nu} - \Omega_{A\rho\nu}\Omega_{P\pi\mu})]N_{P}^{i}.$$

- a) Composing this equation with  $H_{ij}B^j_{\beta}$  one obtains the equation of the form (3.5), that is the  $1^{st}$  Codazzi of the  $3^{rd}$  kind.
- b) Composing (3.7) with  $H_{ij}N_B^j$ , we obtain the 2<sup>nd</sup> Codazzi equation of the  $3^{rd}$  kind in the above cited structure:

$$(3.8) R_{Aij\mu\nu}N_A^iN_B^j = \overline{R}_{AB\mu\nu} + e_A e_B h^{\pi\rho} (\Omega_{A\rho\mu}\Omega_{B\pi\nu} - \Omega_{A\rho\nu}\Omega_{B\pi\mu}).$$

From exposed, the following theorems are valid.

**Theorem 3.1.** The 3<sup>rd</sup> kind integrability conditions of derivational equations (2.1,3) for  $(X_M \subset GR_N, \text{ with the structure } (X_M \subset GR_N, \nabla_{\theta}, \theta \in \{3,4\}), \text{ where }$ the connection  $\nabla$  is defined in (2.6), are given:

- for tangents  $B^i_{\alpha}$  by equation (3.3), for normals  $N^i_A$  by equation (3.7).

**Theorem 3.2.** In the same structure (from the previous theorem) the Gauss equation of the 3<sup>rd</sup> kind for  $X_M \subset GR_N$  is given in (3.4), the 1<sup>st</sup> Codazzi equation of the  $3^{rd}$  kind by (3.5), and the  $2^{nd}$  Codazzi equation of the  $3^{rd}$  kind by (3.8).

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