



## On the Existence of Global Solutions for a Nonlinear Klein-Gordon Equation

Necat Polat<sup>a</sup>, Hatice Taskesen<sup>b</sup>

<sup>a</sup>Department of Mathematics, Dicle University, 21280, Diyarbakir

<sup>b</sup>Department of Statistics, Yuzuncu Yil University, 65080, Van

**Abstract.** The aim of this work is to study the global existence of solutions for the Cauchy problem of a Klein-Gordon equation with high energy initial data. The proof relies on constructing a new functional, which includes both the initial displacement and the initial velocity: with sign preserving property of the new functional we show the existence of global weak solutions.

### 1. Introduction

The nonlinear Klein-Gordon equation with quadratic nonlinearity is

$$u_{tt} - u_{xx} + \alpha u - \beta u^2 = 0, \quad (1)$$

where  $\alpha$ , and  $\beta \neq 0$ . Eq. (1) arises in many scientific applications such as solid state physics, nonlinear optics and quantum field theory. The Klein-Gordon equation is the first relativistic equation in quantum mechanics for the wave function of a particle with zero spin. It was proposed as a relativistic generalization of the Schrödinger equation and was investigated in many papers [1, 2, 4–6, 9, 12, 15, 23, 26].

The goal of the present paper is to investigate the existence of global solutions for the Cauchy problem of the Klein-Gordon equation with dissipation

$$u_{tt} - \Delta u + u + u_t = |u|^{p-1} u, \quad x \in \mathbb{R}^n, \quad t > 0, \quad (2)$$

$$u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), \quad x \in \mathbb{R}^n, \quad (3)$$

where  $u_0$  and  $u_1$  are the initial value functions,  $n \geq 2$  and  $1 < p < \frac{n+2}{n-2}$  if  $n \geq 3$ ,  $1 < p < \infty$  if  $n = 2$ . Evolution equations with dissipation are studied from various aspects in many papers [3, 13, 16, 17, 19].

In the present paper, we investigate the existence of global solutions by using the potential well method [18]. Sattinger [18] investigated global existence of the initial-boundary value problem of the following nonlinear hyperbolic equation

$$u_{tt} - \nabla^2 u + f(x, u) = 0$$

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*Email addresses:* npolat@dicle.edu.tr (Necat Polat), haticetaskesen@yyu.edu.tr (Hatice Taskesen)

in the case of initial energy less than the potential well depth  $d$ . Then this result extended to the total energy of the initial data is less than or equal to  $d$  [24]. Very recently in a paper of Kutev et al. [8] it was proved that there exist global solutions when the total energy of initial data is greater than  $d$  and they established the existence of global weak solutions by constructing a functional which include both the initial displacement and the initial velocity. Because they showed numerically that the initial velocity plays a crucial role in the behaviour of the problem. Problem (2), (3) was already treated in the  $E(0) \leq d$  case by Runzhang [25], but the functional  $I(u)$  used in their paper fails to prove the  $E(0) > d$  case. Although a strongly damped nonlinear Klein-Gordon equation is studied in [26] and a blow up result was given for the high energy initial data, i.e.  $E(0) > d$ , the global existence was studied for  $E(0) \leq d$ . In the present paper, we reinvestigate the problem for the case  $E(0) \leq d$ , where we use a standard functional that include only the initial displacement  $u_0$ . Then, we prove that the existence of global solutions for  $E(0) > d$  can not be proved via sign invariance of this functional. A new functional which includes both the initial displacement  $u_0$  and initial velocity  $u_1$  will be constructed for the case of high energy initial data. Functionals depending on  $u_0$  and  $u_1$  are introduced for the first time in [8] and then they were successfully applied for proving the global existence to some Boussinesq-type equations in [20–22].

Throughout this paper  $H^s = H^s(R^n)$  will denote the  $L^2$  Sobolev space on  $R^n$  with norm  $\|f\|_{H^s} = \|(I - \Delta)^{\frac{s}{2}} f\| = \left\| (1 + k^2)^{\frac{s}{2}} \widehat{f} \right\|$ , where  $s$  is a real number,  $I$  is unitary operator. The notation  $\|f\|_p, \|f\|$  and  $\|f\|_\infty$  will be used instead of norms of  $L^p(R^n), L^2(R^n)$  and  $L^\infty(R^n)$ , respectively.

### 2. Global Existence for $E(0) \leq d$

The present section refers to two points. Firstly, we define a functional which includes only the initial displacement, and prove the existence of global solutions for  $E(0) \leq d$  by aid of the sign invariance of this functional. We then show that this functional fails to prove the global existence in the case of  $E(0) > d$ .

Now, let us define

$$E(t) = E(u(t), u_t(t)) = \frac{1}{2} [\|u_t\|^2 + \|\nabla u\|^2 + \|u\|^2] - \frac{1}{p+1} \|u\|_{p+1}^{p+1}, \tag{4}$$

$$E(t) + \int_0^t \|u_\tau\|^2 d\tau = E(0)$$

$$J(u) = \frac{1}{2} (\|\nabla u\|^2 + \|u\|^2) - \frac{1}{p+1} \|u\|_{p+1}^{p+1}, \tag{5}$$

$$I(u) = (\|\nabla u\|^2 + \|u\|^2) - \|u\|_{p+1}^{p+1}, \tag{6}$$

$$d = \inf_{u \in N} J(u), \tag{7}$$

where  $N = \{u \in H^1 \mid I(u) = 0, \|u\|_{H^1} \neq 0\}$ ,  $E(u(t), u_t(t))$  is the total energy,  $J(u)$  is the potential energy and  $d$  is the depth of potential well which can exactly be written in terms of the Sobolev constant as

$$d = \frac{p-1}{2(p+1)} (S_p^{p+1})^{-2/(p-1)}. \tag{8}$$

Here  $S_p$  is the imbedding constant from  $H^1(R^n)$  into  $L^{p+1}(R^n)$  given by

$$S_p = \sup_{u \in H^1} \frac{\|u\|_{p+1}}{\|u\|_{H^1}}.$$

When  $0 < E(0) < d$ , by the sign invariance of (6) one can prove the existence of global solutions of (2), (3). Existence of global solutions was proved by such functionals for problem (2), (3) in [25]. It was proved in

[25] that if  $I(u) > 0$ , then every weak solutions of the problem exist globally, and if  $I(u) < 0$ , then every weak solutions of the problem blow up in finite time.

For  $\sigma > -\frac{p-1}{2}$ , define

$$\begin{aligned} I_\sigma(u) &= (1 - \sigma) (\|\nabla u\|^2 + \|u\|^2) - \|u\|_{p+1}^{p+1} \\ &= I(u) - \sigma (\|\nabla u\|^2 + \|u\|^2). \end{aligned}$$

Then  $D_\sigma$  and  $N_\sigma$  are defined by

$$D_\sigma = \inf_{u \in N_\sigma} J(u), \quad N_\sigma = \{u \in H^1 : I_\sigma(u) = 0, \|u\|_{H^1} \neq 0\}.$$

Obviously, taking  $\sigma = 0$ ,  $I_\sigma$  corresponds to the functional  $I(u)$ . Moreover, if  $\sigma < -\frac{p-1}{2}$  then  $D_\sigma < 0$ . In this case for  $E(0) = D_\sigma < 0$ , all weak solutions of (2), (3) blow-up in a finite time.

For  $\sigma \in (-\frac{p-1}{2}, 1)$ , we have the following lemmas.

**Lemma 2.1.** Assume that  $u \in H^1(\mathbb{R}^n)$ . If  $I_\sigma(u) < 0$ , then  $\|u\|_{H^1} > \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{1/(p-1)}$ . If  $I_\sigma(u) = 0$ , then  $\|u\|_{H^1} \geq \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{1/(p-1)}$  or  $\|u\|_{H^1} = 0$ .

*Proof.* First, since  $I_\sigma(u) < 0$ , we have  $\|u\|_{H^1} \neq 0$ . Hence, from

$$(1 - \sigma) \|u\|_{H^1}^2 < \|u\|_{p+1}^{p+1} \leq S_p^{p+1} \|u\|_{H^1}^{p+1}, \tag{9}$$

we have  $\|u\|_{H^1} > \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{1/(p-1)}$ .

If  $\|u\|_{H^1} = 0$ , then  $I_\sigma(u) = 0$ . If  $I_\sigma(u) = 0$  and  $\|u\|_{H^1} \neq 0$ , then from

$$(1 - \sigma) \|u\|_{H^1}^2 = \|u\|_{p+1}^{p+1} \leq S_p^{p+1} \|u\|_{H^1}^{p+1}$$

it follows that  $\|u\|_{H^1} \geq \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{1/(p-1)}$ .  $\square$

**Lemma 2.2.** If  $\|u\|_{H^1} < \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{1/(p-1)}$ , then  $I_\sigma(u) > 0$ .

*Proof.* By  $\|u\|_{H^1} < \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{1/(p-1)}$ , we obtain

$$\|u\|_{p+1}^{p+1} \leq S_p^{p+1} \|u\|_{H^1}^{p+1} < (1 - \sigma) \|u\|_{H^1}^2$$

from which follows  $I_\sigma(u) > 0$ .  $\square$

**Theorem 2.3.** Let  $D_\sigma$  be defined as above. Then for  $\sigma > -\frac{p-1}{2}$ , we have

$$D_\sigma = \frac{p-1+2\sigma}{2(p+1)} \left(\frac{|1-\sigma|}{S_p^{p+1}}\right)^{2/(p-1)}.$$

If we write  $D_\sigma$  in terms of  $d$ , we obtain

$$D_\sigma = \frac{p-1+2\sigma}{2(p+1)} |1-\sigma|^{2/(p-1)} \frac{2(p+1)}{p-1} d. \tag{10}$$

*Proof.* If  $u \in N_\sigma$ , we have by Lemma 2.1 that  $\|u\|_{H^1} \geq \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{1/(p-1)}$ . In the proof of Lemma 2.1 the inequality (9) is an equality iff  $u$  is a minimizer of the imbedding  $H^1$  into  $L^{p+1}$ . Since  $\|u\|_{p+1} = S_p \|u\|_{H^1}$  is attained only for  $\tilde{u} = \left(\cosh\left(\frac{p-1}{2}x\right)\right)^{-\frac{2}{p-1}}$  for  $n = 1$  [11], and for the ground state solution of (2), (3) for  $n > 1$  [14] and it has constant sign, we have

$$\inf_{u \in N_\sigma} \|u\|_{H^1} = \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{1/(p-1)}.$$

Hence from

$$\begin{aligned} \inf_{u \in N_\sigma} J(u) &= \inf_{u \in N_\sigma} \left(\frac{1}{2} \|u\|_{H^1}^2 - \frac{1}{p+1} \|u\|_{p+1}^{p+1}\right) \\ &= \inf_{u \in N_\sigma} \left[\left(\frac{1}{2} - \frac{(1-\sigma)}{p+1}\right) \|u\|_{H^1}^2 + \frac{1}{p+1} I_\sigma(u)\right] \\ &= \left(\frac{1}{2} - \frac{(1-\sigma)}{p+1}\right) \inf_{u \in N_\sigma} \|u\|_{H^1}^2, \end{aligned}$$

and by definition of  $D_\sigma$  we obtain  $D_\sigma = \frac{p-1+2\sigma}{2(p+1)} \left(\frac{1-\sigma}{S_p^{p+1}}\right)^{2/(p-1)}$ .  $\square$

We can also state the following properties of  $D_\sigma$ , which can be proved easily.

- i)  $D_\sigma$  is strictly increasing on  $\sigma \in \left(-\frac{p-1}{2}, 0\right) \cup (1, \infty)$  and strictly decreasing on  $(0, 1)$ .
- ii)  $\lim_{\sigma \rightarrow 1} D_\sigma = 0$ , and  $D_{\sigma_0} = 0$ , where  $\sigma_0 = -\frac{p-1}{2}$ .

The following theorems show the invariance of  $I_\sigma$  under the flow of (2), (3) for  $0 < E(0) < d$  and  $E(0) = d$ , respectively, and can be proved by contradiction as in [20].

**Theorem 2.4.** Assume that  $u_0 \in H^1(R^n)$ ,  $u_1 \in L^2(R^n)$ . Let  $0 < E(0) < d$ . Then the sign of  $I_\sigma$  is invariant under the flow of (2), (3) for  $\sigma \in (\sigma_1, \sigma_2]$ , where  $\sigma_1$  and  $\sigma_2$  are the corresponding minimal negative and minimal positive roots of equation  $D_\sigma = E(0)$ .

**Theorem 2.5.** Let all the assumptions of Theorem 2.4 hold and that  $E(0) = d$ . Then the sign of  $I_0$  (recall that when  $E(0) = d$ , we have  $\sigma_1 = \sigma_2 = 0$ ) is invariant with respect to (2), (3) for every  $t \in [0, \infty)$ .

Now, we give a lemma for  $\sigma > 1$ , which states similar results to Lemmas 2.1, 2.2, and can be proved similarly.

**Lemma 2.6.** Assume that  $u \in H^1(R^n)$ . Let  $\sigma > 1$ . If  $I_\sigma(u) > 0$ , then  $\|u\|_{H^1} > s(\sigma)$ . If  $I_\sigma(u) = 0$ , then  $\|u\|_{H^1} \geq s(\sigma)$  or  $\|u\|_{H^1} = 0$ , where  $s(\sigma) = \left(\frac{\sigma-1}{S_p^{p+1}}\right)^{1/(p-1)}$ . Moreover, if  $\|u\|_{H^1} < s(\sigma)$ , then  $I_\sigma(u) \leq 0$  and  $I_\sigma(u) = 0$  if and only if  $\|u\|_{H^1} = 0$ .

**Theorem 2.7.** Assume that  $u_0 \in H^1(R^n)$ ,  $u_1 \in L^2(R^n)$ . If  $E(0) > 0$ , then  $I_\sigma(u(t)) \leq 0$  for every  $t > 0$  and  $\sigma \geq \sigma_m$ , where  $\sigma_m$  is the maximal positive root of  $D_\sigma = E(0)$ .

*Proof.* We give the proof of the theorem for  $\sigma = \sigma_m$  and  $\sigma > \sigma_m$  separately. First, we prove the theorem for  $\sigma = \sigma_m$ . By contradiction, assume that there exists some  $t' > 0$  such that  $I_{\sigma_m}(u(t')) > 0$ . By Lemma 2.1, we have  $\|u\|_{H^1} > 0$  and there exists a value  $\sigma > \sigma_m$  such that  $I_\sigma(u(t')) = 0$ . Then, by (4),  $D_{\sigma_m} = E(0) \geq J(u(t')) \geq \inf_{u \in N_\sigma} J(u) = D_\sigma$ . By definition of  $D_\sigma$ , for  $\sigma > \sigma_m > 1$  we have  $D_\sigma > D_{\sigma_m}$ . A contradiction occurs, which proves the theorem for  $\sigma = \sigma_m$ . For  $\sigma \geq \sigma_m$ ,  $I_{\sigma_m}(u(t)) \geq I_\sigma(u(t))$  implies that the theorem is true for every  $\sigma \geq \sigma_m$ .  $\square$

The following Corollary gives a more precise result for subcritical initial energy.

**Corollary 2.8.** *Suppose  $u_0 \in H^1(R^n)$ ,  $u_1 \in L^2(R^n)$ . Let  $0 < E(0) < d$  and  $I_0(u_0) > 0$ . Then,*

$$0 < I_0(u(t)) \leq \sigma_m \|u\|_{H^1}^2 \tag{11}$$

for every  $t > 0$ .

*Proof.* We know that for  $I_0(u(t)) > 0$ , the solution  $u(x, t)$  of problem (2), (3) is globally defined. Since  $E(0) = D_{\sigma_m}$  for some  $\sigma_m > 1$  then by Theorem 2.7 we have  $I_{\sigma_m}(u(t)) \leq 0$  for every  $t \in [0, \infty)$ . Thus we get the inequality (11) from below and from above.  $\square$

**Remark 2.9.** *We tried to characterize the behavior of solutions for  $E(0) > d$  in terms of initial displacement. We constituted the new functional  $I_\sigma(u)$  and proved the sign invariance of  $I_\sigma(u)$  for  $0 < E(0) < d$  and  $E(0) = d$ . But the case  $E(0) > d$  is still an open question, because from Theorem 2.7, we concluded that in this case  $I_\sigma(u)$  is always non-positive.*

### 3. Main Results

We will introduce our new functional which will be used for global existence of solutions with high energy initial data.

$$\tilde{M}(v, \omega) = (\|\nabla v\|^2 + \|v\|^2) - \|v\|_{p+1}^{p+1} - (\omega, \omega) \tag{12}$$

for every  $v \in H^1$  and  $\omega \in L^2$ . For simplicity we denote

$$M(u, t) = \tilde{M}(u(\cdot, t), u_t(\cdot, t)).$$

The sign invariance of this new functional can be stated as follows.

**Theorem 3.1.** *Let  $u_0 \in H^1(R^n)$ ,  $u_1 \in L^2(R^n)$  and  $E(0) > 0$ . For  $\sigma > \sigma_m$ , assume that*

$$(u_1, u_0) + \frac{1}{2} \|u_0\|^2 + \frac{(p+1)\sigma}{p-1+(p+3)\sigma} E(0) \leq 0. \tag{13}$$

*If  $M(u, 0)$  is positive, then  $M(u, t)$  is positive for every  $t \in [0, \infty)$ .*

*Proof.* [Proof] We prove the theorem by contradiction. Let us define

$$\theta(t) = \|u\|^2 + \int_0^t \|u\|^2 d\tau.$$

Then

$$\theta'(t) = 2(u_t, u) + \|u\|^2,$$

$$\begin{aligned} \theta''(t) &= 2\|u_t\|^2 + 2(u_{tt}, u) + 2(u_t, u) \\ &= 2\|u_t\|^2 + 2\left[\|u\|_{p+1}^{p+1} - \|\nabla u\|^2 - \|u\|^2 - (u_t, u)\right] + 2(u_t, u) \\ &= -2M(u, t). \end{aligned}$$

To get a contradiction, let us assume that there exists some  $t' > 0$  such that  $M(u, t') = 0$ . Since  $\theta''(t) < 0$ , we conclude that  $\theta'(t)$  is strictly decreasing on  $[0, t')$ . Moreover, (13) implies  $\theta'(0) < 0$  and therefore  $\theta'(t) < 0$  in

$[0, t']$ , from which follows that  $\theta(t)$  is strictly decreasing on  $[0, t']$ . By the energy identity and  $M(u, t') = 0$ , we have

$$\begin{aligned} E(0) &\geq \frac{1}{2} \|u_t(t')\|^2 + \frac{p-1}{2(p+1)} (\|\nabla u\|^2 + \|u\|^2) + \frac{1}{p+1} I(u(t')) \\ &= \left(\frac{1}{2} + \frac{1}{p+1}\right) \|u_t(t')\|^2 + \frac{p-1}{2(p+1)} (\|\nabla u\|^2 + \|u\|^2). \end{aligned} \tag{14}$$

Theorem 2.7 and  $M(u, t') = 0$  yield

$$(\|\nabla u\|^2 + \|u\|^2) \geq \sigma_m^{-1} I_0(u(t')) \geq \sigma^{-1} \|u_t(t')\|^2.$$

The use of this inequality in (14) gives

$$\begin{aligned} E(0) &\geq \left(\frac{1}{2} + \frac{1}{p+1} + \frac{p-1}{2(p+1)\sigma}\right) \|u_t(t')\|^2 \\ &= \frac{(p+3)\sigma + p-1}{2(p+1)\sigma} [\|(u_t(t') + u(t'))\|^2 \\ &\quad - 2(u_t(t'), u(t')) - \|u(t')\|^2]. \end{aligned}$$

From the monotonicity of  $\theta(t)$  and  $\theta'(t)$ , we get

$$E(0) > \frac{(p+3)\sigma + p-1}{(p+1)\sigma} \left[-(u_1, u_0) - \frac{1}{2} \|u_0\|^2\right]$$

which contradicts with (13). Thus the proof is completed.  $\square$

**Theorem 3.2.** *Let  $1 < p < \infty$  for  $n = 2$ ;  $1 < p < \frac{n+2}{n-2}$  for  $n \geq 3$  and  $u_0 \in H^1(\mathbb{R}^n)$ ,  $u_1 \in L^2(\mathbb{R}^n)$ . Suppose that  $E(0) > 0$ ,  $M(u, 0) > 0$  and (13) holds for some  $\sigma > \sigma_m$ . Then, the weak solution of problem (2),(3) is globally defined for every  $t \in [0, \infty)$ .*

*Proof.* [Proof]The proof of this theorem follows from adding some arguments to the local existence result of Proposition 1.1 of [25].  $M(u, 0) > 0$  implies from the sign preserving property of  $M(u, t)$  that  $M(u, t) > 0$ , thereby  $I_0(u) > 0$  for every  $t > 0$ . From energy identity, we have

$$\begin{aligned} E(0) &\geq \frac{1}{2} \|u_t\|^2 + \frac{p-1}{2(p+1)} (\|\nabla u\|^2 + \|u\|^2) + \frac{1}{p+1} I(u) \\ &\geq \frac{1}{2} \|u_t\|^2 + \frac{p-1}{2(p+1)} (\|\nabla u\|^2 + \|u\|^2). \end{aligned}$$

Therefore  $\|u\|_{H^1}$  and  $\|u_t\|_{L^2}$  are bounded for every  $t > 0$ . The previously mentioned local existence theory completes the proof.  $\square$

### References

- [1] J.D. Avrin, Convergence Properties of the Strongly Damped Nonlinear Klein-Gordon Equation, J. Differential Equations 67 (1987) 243-255.
- [2] T. Cazenave, Uniform Estimates for Solutions of Nonlinear Klein-Gordon Equations, J. Func. Anal.60 (1985) 36-55.
- [3] F. Gazzola, M. Squassina, Global solutions and finite time blow up for damped semilinear wave equations, Ann. I. H. Poincaré – AN 23 (2006) 185–207.
- [4] J. Ginibre, G. Velo, The global Cauchy problem for the nonlinear Klein Gordon equation, Math. Z. 189 (4) (1985) 487-505.
- [5] J. Ginibre, G. Velo, The global Cauchy problem for the nonlinear Klein Gordon equation II, Ann. Inst. H. Poincaré Anal. Non Linéaire 6 (1) (1989) 15-35.
- [6] A.I. Komech, E.A. Kopylova, Weighted energy decay for 3D Klein–Gordon equation, J. Differential Equations 248 (2010) 501-520.

- [7] N. Kutev, N. Kolkovska, M. Dimova, C.I. Christov, Theoretical and numerical aspects for global existence and blow up for the solutions to Boussinesq paradigm equation, *AIP Conf. Proc.* 1404 (2011) 68–76.
- [8] N. Kutev, N. Kolkovska, M. Dimova, Global existence of Cauchy problem for Boussinesq paradigm equation, *Comput. Math. Appl.*, 65 (2013) 500-511.
- [9] C. Laurent, On stabilization and control for the critical Klein–Gordon equation on a 3-D compact manifold, *J. Func. Anal.* 260 (2011) 1304-1368.
- [10] H. A. Levine, Instability and nonexistence of global solutions of nonlinear wave equation of the form  $Pu = Au_{tt} + F(u)$ , *TAMS* 192 (1974) 1–21.
- [11] E. Lieb, Sharp constants in the Hardy-Littlewood-Sobolev and related inequalities, *Annals of Math.* 118 (1983) 349–374.
- [12] Y. Liu, M. Ohta, G. Todorova, Strong instability of solitary waves for nonlinear Klein–Gordon equations and generalized Boussinesq equations, *Ann. I. H. Poincaré – AN* 24 (2007) 539–548.
- [13] L. Liu, M. Wang, Global existence and blow-up of solutions for some hyperbolic systems with damping and source terms, *Nonlinear Analysis* 64 (2006) 69 – 91.
- [14] Y. Liu, Instability and blow-up of solutions to a generalized Boussinesq equation, *SIAM J. Math. Anal.* 26 (1995) 1527-1546.
- [15] M. Nakamura, The Cauchy problem for semi-linear Klein–Gordon equations in de Sitter spacetime, *J. Math. Anal. Appl.* 410 (2014) 445-454.
- [16] E. Pişkin, N. Polat, Uniform decay and blow up of solutions for a system of nonlinear higher-order Kirchhoff-type equations with damping and source, *Contemp. Anal. Appl. Math.* 1(2) (2013) 181-199.
- [17] N. Polat, A. Ertaş, Existence and blow-up of solution of Cauchy problem for the generalized damped multidimensional Boussinesq equation, *J. Math. Anal. Appl.* 349(1) (2009) 10-20.
- [18] D.H. Sattinger, On global solution of nonlinear hyperbolic equations, *Arch. Rational Mech. Anal.* 30 (1968) 148-172.
- [19] H. Soltanov, A note on the Goursat problem for a multidimensional hyperbolic equation, *Contemp. Anal. Appl. Math.* 1(2) (2013) 98-106.
- [20] H. Taskesen, N. Polat, A. Ertaş, On global solutions for the Cauchy problem of a Boussinesq-type equation, *Abst. Appl. Anal.* 2012 (2012) 10 pages.
- [21] H. Taskesen, N. Polat, Existence of global solutions for a multidimensional Boussinesq-type equation with supercritical initial energy, in: A. Ashyralyev, A. Lukashov (Eds.), *First international conference on analysis and applied mathematics: ICAAM 2012*, volume 1407 of *AIP Conf. Proc.*, pp.159-162.
- [22] H. Taskesen, N. Polat, Global existence for a double dispersive sixth order Boussinesq equation, *Contemp. Anal. Appl. Math.*, 1(1) (2013) 60-69.
- [23] B. Wang, Scattering of solutions for critical and subcritical nonlinear Klein-Gordon equations in  $H^s$ , *Discrete and Continuous Dynamical Systems* 5(4) (1999) 753-763.
- [24] S. Wang, H. Wang, Global solution for a generalized Boussinesq equation, *Appl. Math. Comput.* 204 (2008) 130–136.
- [25] R. Xu, Global existence, blow up and asymptotic behaviour of solutions for nonlinear Klein–Gordon equation with dissipative term, *Math. Meth. Appl. Sci.* 33 (2010) 831-844.
- [26] R. Xu, Y. Ding, Global solutions and finite time blow up for damped Klein-Gordon equation, *Acta Mathematica Scientia* 33B(3) (2013) 643-652.