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About Almost Geodesic Curves

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Abstract. We determine in \mathbb{R}^n the form of curves C for which also any image under an (n-1)-dimensional algebraic torus is an almost geodesic with respect to an affine connection ∇ with constant coefficients and calculate the components of ∇ .

1. Introduction

This paper is a result following on from D. Betten researchs [4] and our papers [1, 8–10].

The geodesics and almost geodesics play an important role in differential geometry. For this reason many geometricians study almost geodesic mappings (see [3], [14], [17]). In [12], [13] almost geodesic curves were considered in generalized Riemannian and Kählerian spaces. E. Beltrami [6] has shown that a differentiable curve is a local geodesic with respect to an affine connection ∇ precisely if it is a solution of an Abelian differential equation with coefficients which are functions of the components of ∇ . The investigation with systems of lines of 2-dimensional topological geometries was started in [15]. The explicit calculation of the form of curves C in the n-dimensional real space \mathbb{R}^n which are geodesics or almost geodesics with respect to an affine connection ∇ is not achievable even in the case if the components Γ^h_{ij} of ∇ are constant. But we did it. In [2] the geodesics and special case of almost geodesics were considered. We supposed that with C also all images of C under a real (n-1)-dimensional algebraic torus are also geodesics, respectively almost geodesics. This implies that the determination of C becomes an algebraic problem (a problem of polynomial identities). Our model allows you to look at known things globally. In this paper we continue to study almost geodesic curves [7], [16] and here we will consider other case.

We consider a curve C homeomorphic to $\mathbb R$ which is a closed subset of $\mathbb R^n$ and has the form

$$C = (t, f_2(t), \dots, f_n(t)), t \in \mathbb{R}, \tag{1}$$

where $f_i(t)$: $\mathbb{R} \to \mathbb{R}$, i = 2, ..., n, are three times differentiable non-constant functions. The system

$$\mathfrak{X}(C) = \{(t + c_1, b_2 f_2(t) + c_2, \dots, b_n f_n(t) + c_n), t \in \mathbb{R}\}, \text{ where } b_i \neq 0, c_i \in \mathbb{R},$$

is a set of imagines of *C*.

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If every curve of $\mathfrak{X}(C)$ is a geodesic with respect to an affine connection ∇ with constant coefficients $\Gamma^h_{ij'}$ then the derivatives $f'_i(t)$ of the functions $f_i(t)$ are solutions of the first order linear ordinary differential equations. If every curve of $\mathfrak{X}(C)$ is an almost geodesic with respect to ∇ , then the derivatives $f'_i(t)$ are solutions of harmonic oscillator equations. If $\mathfrak{X}(C)$ consists of Euclidean lines which are geodesics with respect to ∇ , then at the most Γ^1_{11} may be different from 0. In contrast to this if $\mathfrak{X}(C)$ consists of Euclidean lines then there is huge quantity of non-trivial connections ∇ such that the lines of $\mathfrak{X}(C)$ are almost geodesic with respect to ∇ .

Since we apply results of differential geometry only for the n-dimensional space \mathbb{R}^n , where global coordinates exist and the components Γ^h_{ij} , $h, i, j \in \{1, 2, ..., n\}$, of any affine connection ∇ can be written in unique way in these coordinates.

Remark 1.1. If coefficients Γ_{ii}^h of an affine connection ∇ are constants then there exist groups of affine movements.

Remark 1.2. It is possible to apply our model for other spaces, because a geodesic and an almost geodesic can be defined in other spaces in the same manner as in \mathbb{R}^n [11].

2. Almost geodesic curves

Let

$$\ell = (t + c_1, b_2, f_2(t) + c_2, \dots, b_n, f_n(t) + c_n), t \in \mathbb{R},$$

be a curve of $\mathfrak{X}(C)$. Then

$$\dot{\ell} = (1, b_2 f_2'(t), \dots, b_n f_n'(t)), \quad \ddot{\ell} = (0, b_2 f_2''(t), \dots, b_n f_n''(t)).$$

By an *almost geodesic* of an affine connection ∇ we mean a piecewise C^3 -curve $\gamma: I \to \mathbb{R}^n$ satisfying

$$\nabla_{\dot{\gamma}}(\nabla_{\dot{\gamma}}\dot{\gamma}) = \rho \cdot \dot{\gamma} + \sigma \cdot \nabla_{\dot{\gamma}}\dot{\gamma},$$

where ϱ , σ : $I \to \mathbb{R}$ are continuous functions, $I \subset \mathbb{R}$ is an open interval (cf. [16, p. 158], [7, p. 456]). Using the components of ∇ the system of differential equations for almost geodesics has the form

$$\ddot{\gamma}^h + \sum_{i,j,k=1}^n (\partial_k \Gamma^h_{ij} + \Gamma^\ell_{ij} \Gamma^h_{\ell k}) \dot{\gamma}^i \dot{\gamma}^j \dot{\gamma}^k + 2 \sum_{i,j=1}^n \Gamma^h_{ij} \dot{\gamma}^i \dot{\gamma}^j + \sum_{i,j=1}^n \Gamma^h_{ij} \dot{\gamma}^i \dot{\gamma}^j = \varrho(t) \cdot \dot{\gamma}^h + \sigma(t) \cdot (\dot{\gamma}^h + \sum_{i,j=1}^n \Gamma^h_{ij} \dot{\gamma}^i \dot{\gamma}^j). \tag{2}$$

A curve ℓ of $\mathfrak{X}(C)$ is an almost geodesic with respect to a connection ∇ with constant coefficients $\{\Gamma_{ij}^h\}$ if and only if according to (2) we have

$$\ddot{\ell}^h + \sum_{i,j,k=1}^n \Gamma^{\ell}_{ij} \Gamma^h_{\ell k} \, \dot{\ell}^i \dot{\ell}^j \dot{\ell}^k + 2 \sum_{i,j=1}^n \Gamma^h_{ij} \ddot{\ell}^i \dot{\ell}^j + \sum_{i,j=1}^n \Gamma^h_{ij} \dot{\ell}^i \ddot{\ell}^j = \varrho(t) \cdot \dot{\ell}^h + \sigma(t) \cdot (\ddot{\ell}^h + \sum_{i,j=1}^n \Gamma^h_{ij} \dot{\ell}^i \dot{\ell}^j). \tag{3}$$

We rewrite the formula (3) for h = 1 and obtain the function $\varrho(t)$. For h = 2, ..., n after substitution $\varrho(t)$ in (3) we get

$$b_{h} f_{h}^{\prime\prime\prime}(t) + \sum_{m=1}^{n} \Gamma_{m1}^{h} \Gamma_{11}^{m} + \sum_{i=2,m=1}^{n} (\Gamma_{mi}^{h} \Gamma_{11}^{m} + \Gamma_{m1}^{h} \Gamma_{i1}^{m} + \Gamma_{m1}^{h} \Gamma_{1i}^{m}) b_{i} f_{i}^{\prime}(t) +$$

$$\sum_{i,j=2,m=1}^{n} (\Gamma_{mi}^{h} \Gamma_{j1}^{m} + \Gamma_{mi}^{h} \Gamma_{1j}^{m} + \Gamma_{m1}^{h} \Gamma_{ij}^{m}) b_{i} b_{j} f_{i}^{\prime}(t) f_{j}^{\prime}(t) +$$

$$\sum_{i,j,k=2,m=1}^{n} \Gamma_{mi}^{h} \Gamma_{jk}^{m} b_{i} b_{j} b_{k} f_{i}^{\prime}(t) f_{j}^{\prime}(t) f_{k}^{\prime}(t) +$$

$$\sum_{i=2}^{n} (2\Gamma_{i1}^{h} + \Gamma_{1i}^{h}) b_{i} f_{i}^{"}(t) + \sum_{i,j=2}^{n} (2\Gamma_{ij}^{h} + \Gamma_{ji}^{h}) b_{i} b_{j} f_{i}^{"}(t) f_{j}^{'}(t) - b_{i} f_{h}^{'}(t) \left(\sum_{m=1}^{n} \Gamma_{m1}^{1} \Gamma_{11}^{m} + \sum_{i=2,m=1}^{n} (\Gamma_{mi}^{1} \Gamma_{11}^{m} + \Gamma_{m1}^{1} \Gamma_{i1}^{m} + \Gamma_{m1}^{1} \Gamma_{ij}^{m}) b_{i} b_{j} f_{i}^{'}(t) + \sum_{i,j=2,m=1}^{n} (\Gamma_{mi}^{1} \Gamma_{j1}^{m} + \Gamma_{mi}^{1} \Gamma_{1j}^{m} + \Gamma_{m1}^{1} \Gamma_{ij}^{m}) b_{i} b_{j} f_{i}^{'}(t) f_{j}^{'}(t) + \sum_{i,j=2,m=1}^{n} \Gamma_{mi}^{1} \Gamma_{jk}^{m} b_{i} b_{j} b_{k} f_{i}^{'}(t) f_{j}^{'}(t) f_{k}^{'}(t) + \sum_{i,j=2}^{n} (2\Gamma_{i1}^{1} + \Gamma_{1i}^{1}) b_{i} f_{i}^{"}(t) + \sum_{i,j=2}^{n} (2\Gamma_{ij}^{1} + \Gamma_{ji}^{1}) b_{i} b_{j} f_{i}^{"}(t) f_{j}^{'}(t) + \sum_{i,j=2}^{n} (\Gamma_{i1}^{h} - \Gamma_{1i}^{1} - \Gamma_{1i}^{1}) b_{i} f_{i}^{'}(t) + \sum_{i,j=2}^{n} (\Gamma_{ij}^{h} - \Gamma_{ij}^{1}) b_{i} b_{j} f_{i}^{'}(t) f_{j}^{'}(t) \right).$$

$$(4)$$

One can determine σ only if not all coefficients in (4) are zero. In [2] we treated the case that for $h \ge 2$ one has

$$\Gamma^{h}_{i1} + \Gamma^{h}_{1i} = \Gamma^{1}_{i1} + \Gamma^{1}_{1i}, \ i \ge 1, \ \text{ and } \ \Gamma^{h}_{ij} = \Gamma^{1}_{ij} \text{ for all } i, j \ge 2.$$

Now let an α and i_0 , j_0 such that for these indices we have

$$\Gamma_{11}^{\alpha} \neq \Gamma_{11}^{1} \text{ or } \Gamma_{i_01}^{\alpha} + \Gamma_{1i_0}^{\alpha} \neq \Gamma_{i_01}^{1} + \Gamma_{1i_0}^{1}, \text{ or } \Gamma_{i_0j_0}^{\alpha} \neq \Gamma_{i_0j_0}^{1}, i_0, j_0 \ge 2.$$
 (5)

In this case the coefficient of σ is not identically zero, and we can compute σ . Putting the expression of σ into relation (4) we obtain

$$\left((\Gamma_{11}^{\alpha} - \Gamma_{11}^{1}) + 2(\Gamma_{1i_{0}}^{\alpha} + \Gamma_{i_{0}1}^{\alpha} - \Gamma_{1i_{0}}^{1} - \Gamma_{i_{0}1}^{1})f_{i_{0}}'b_{i_{0}} + (\Gamma_{i_{0}j_{0}}^{\alpha} - \Gamma_{i_{0}j_{0}}^{1})f_{i_{0}}'f_{j_{0}}'b_{i_{0}}b_{j_{0}}\right) \cdot \left(T_{h111} + \left(f_{h}''' - T_{1111}f_{h}'\right)b_{h} + \sum_{i=2}^{n} \left(S_{h11i}f_{i}' + (2\Gamma_{i1}^{h} + \Gamma_{1i}^{h})f_{i}''\right)b_{i} - \sum_{i=2}^{n} \left(S_{1i11}f_{i}' + (2\Gamma_{i1}^{1} + \Gamma_{1i}^{1})f_{i}''\right)f_{h}'b_{h}b_{i} + \sum_{i,j=2}^{n} \left(S_{hij1}f_{i}'f_{j}' + (2\Gamma_{ij}^{h} + \Gamma_{ji}^{h})f_{i}''f_{j}'\right)b_{i}b_{j} - \sum_{i,j=2}^{n} \left(S_{11ij}f_{i}'f_{j}' + (2\Gamma_{ij}^{1} + \Gamma_{ji}^{1})f_{i}''f_{j}'\right)f_{h}'b_{h}b_{i}b_{j} + \sum_{i,j=2}^{n} T_{hijk}f_{i}'f_{j}'f_{k}'b_{i}b_{j}b_{k} - \sum_{i,j,k=2}^{n} T_{1ijk}f_{h}'f_{i}'f_{j}'f_{k}'b_{h}b_{i}b_{j}b_{k}\right) = \left((\Gamma_{11}^{h} - \Gamma_{11}^{1}) + 2\sum_{i=2}^{n} (\Gamma_{1i}^{h} + \Gamma_{i1}^{h} - \Gamma_{1i}^{1} - \Gamma_{1i}^{1})f_{i}'b_{i} + \sum_{i,j=2}^{n} (\Gamma_{ij}^{h} - \Gamma_{ij}^{1})f_{i}'f_{j}'b_{i}b_{j}\right) \cdot \left(T_{\alpha 111} + \left(f_{\alpha}''' - T_{1111}f_{\alpha}'\right)b_{\alpha} + \left(S_{\alpha 11i_{0}}f_{i_{0}}' + (2\Gamma_{i01}^{\alpha} + \Gamma_{1i_{0}}^{\alpha})f_{i_{0}}''\right)b_{i_{0}} - \left(S_{11i_{0}j_{0}}f_{i_{0}}' + (2\Gamma_{i01}^{1} + \Gamma_{1i_{0}}^{1})f_{i_{0}}''\right)f_{\alpha}'b_{h}b_{i_{0}} + \left(S_{\alpha i_{0}j_{0}}f_{i_{0}}' + (2\Gamma_{i_{0}}^{\alpha} + \Gamma_{j_{0}i_{0}}^{\alpha})f_{i_{0}}''\right)f_{j_{0}}'b_{i_{0}}b_{j_{0}} - \left(S_{11i_{0}j_{0}}f_{i_{0}}' + (2\Gamma_{i_{0}}^{1} + \Gamma_{1i_{0}}^{1})f_{i_{0}}''\right)f_{\alpha}'f_{j_{0}}b_{a}b_{i_{0}}b_{j_{0}} + \left(S_{\alpha i_{0}j_{0}}f_{i_{0}}' + (2\Gamma_{i_{0}}^{\alpha} + \Gamma_{j_{0}i_{0}}^{1})f_{i_{0}}''\right)f_{\alpha}'f_{j_{0}}b_{a}b_{i_{0}}b_{j_{0}}\right) + \sum_{k=2}^{n} T_{\alpha i_{0}j_{0}}f_{i_{0}}'f_{j_{0}}'f_{k}'b_{i_{0}}b_{j_{0}}b_{k} - \sum_{k=2}^{n} T_{1i_{0}j_{0}}f_{\alpha}'f_{i_{0}}'f_{j_{0}}'f_{k}'b_{a}b_{i_{0}}b_{j_{0}}b_{k}\right),$$
(6)

where

$$S_{ABCD} \stackrel{def}{=} \sum_{m=1}^{n} \left(\Gamma_{mD}^{A} \Gamma_{BC}^{m} + \Gamma_{mB}^{A} (\Gamma_{DC}^{m} + \Gamma_{CD}^{m}) \right),$$
$$T_{ABCD} \stackrel{def}{=} \sum_{m=1}^{n} \Gamma_{mB}^{A} \Gamma_{CD}^{m}.$$

If n = 2, then $h = \alpha = 2$ and from (6) we obtain that any plane curve of the system $\mathfrak{X}(C)$, where C has the form (1), is an almost geodesic if the affine connection Γ_{ij}^h satisfies the conditions (5). Hence we assume $n \ge 3$.

Now we consider the first case, when

$$\Gamma_{11}^h = \Gamma_{11}^1 \text{ for all } 2 \le h \le n \text{ and } \Gamma_{i1}^h + \Gamma_{1i}^h = \Gamma_{i1}^1 + \Gamma_{1i}^1 \text{ for all } 2 \le i \le n,$$
 (7)

but there exists an α and i_0 , j_0 such that

$$\Gamma^{\alpha}_{i_0 i_0} \neq \Gamma^1_{i_0 i_0}. \tag{8}$$

Writing a system of equations and conditions which follow from (6) and using linear independence functions we get differential equations. Integrating them (see. [5]) we obtain the following

Theorem 2.1. Let C be a curve of the form (1) and ∇ be a connection with constant coefficients $\{\Gamma_{ij}^h\}$ satisfying relations (7), (8).

Then any curve ℓ of $\mathfrak{X}(C)$ is almost geodesic with respect to ∇ if and only if ℓ is represented by the functions f_h , f_{α} , f_{i_0} having the following forms

- * $f_h(t) = \hat{C}_h e^{\lambda_1^h t} + \hat{D}_h e^{\lambda_2^h t}$, where \hat{C}_h , $\hat{D}_h \in \mathbb{R}$ are not both zero and $a_h^2 4c_h > 0$,
- * $f_h(t) = (\tilde{C}_h t + \tilde{D}_h)e^{\frac{-a_h}{2}t}$, where \tilde{C}_h , $\tilde{D}_h \in \mathbb{R}$ are not both zero and $a_h^2 4c_h = 0$,
- * $f_h(t) = e^{-a_h t/2} \left(\bar{C}_h \cos \frac{\sqrt{a_h^2 4c_h}}{2} t + \bar{D}_h \sin \frac{\sqrt{a_h^2 4c_h}}{2} t \right)$, where \bar{C}_h , $\bar{D}_h \in \mathbb{R}$ are not both zero and $a_h^2 4c_h < 0$

with

$$a_h = 2\Gamma_{h1}^h + \Gamma_{1h}^h, \quad c_h = S_{h11h} - T_{1111},$$

$$\lambda_1^h = \frac{-a_h - \sqrt{a_h^2 - 4c_h}}{2}, \quad \lambda_2^h = \frac{-a_h + \sqrt{a_h^2 - 4c_h}}{2};$$

- * $f_{\alpha}(t)=C_{\alpha}t^2+D_{\alpha}t+E$, where $C_{\alpha},D_{\alpha},E\in\mathbb{R},\ C_{\alpha},D_{\alpha}$ are not both zero and $\gamma_{\alpha}=0$,
- * $f_{\alpha}(t) = \hat{C}_{\alpha}e^{\sqrt{-\gamma_{\alpha}}t} \hat{D}_{\alpha}e^{-\sqrt{-\gamma_{\alpha}}t}$, where \hat{C}_{α} , $\hat{D}_{\alpha} \in \mathbb{R}$ are not both zero and $\gamma_{\alpha} < 0$,
- * $f_{\alpha}(t) = \hat{C}_{\alpha} \sin(\sqrt{\gamma_{\alpha}} t) \hat{D}_{\alpha} \cos(\sqrt{\gamma_{\alpha}} t)$, where \hat{C}_{α} , $\hat{D}_{\alpha} \in \mathbb{R}$ are not both zero and $\gamma_{\alpha} > 0$

with

$$\gamma_{\alpha} = \frac{(\Gamma_{i_0 j_0}^{\alpha} - \Gamma_{i_0 j_0}^{1})(T_{h\alpha ij} + T_{h\alpha ji} + T_{hi\alpha j} + T_{hij\alpha} + T_{hj\alpha i} + T_{hji\alpha})}{\Gamma_{ij}^{1} + \Gamma_{ji}^{1} - \Gamma_{ij}^{h} - \Gamma_{ji}^{h}} - T_{1111};$$

- * $f_{i_0}(t) = \hat{C}_{i_0} e^{\lambda_1^{i_0} t} + \hat{D}_{i_0} e^{\lambda_2^{i_0} t}$, where \hat{C}_{i_0} , $\hat{D}_{i_0} \in \mathbb{R}$ are not both zero and $a_{i_0}^2 4c_{i_0} > 0$,
- * $f_{i_0}(t) = (\tilde{C}_{i_0}t + \tilde{D}_{i_0})e^{\frac{-a_{i_0}}{2}t}$, where \tilde{C}_{i_0} , $\tilde{D}_{i_0} \in \mathbb{R}$ are not both zero and $a_{i_0}^2 4c_{i_0} = 0$,
- * $f_{i_0}(t) = e^{-a_{i_0}t/2} \left(\bar{C}_{i_0} cos \frac{\sqrt{a_{i_0}^2 4c_{i_0}}}{2}t + \bar{D}_{i_0} sin \frac{\sqrt{a_{i_0}^2 4c_{i_0}}}{2}t\right)$, where \bar{C}_{i_0} , $\bar{D}_{i_0} \in \mathbb{R}$ are not both zero and $a_{i_0}^2 4c_{i_0} < 0$

with

$$a_{i_0} = 2\Gamma^{i_0}_{i_01} + \Gamma^{i_0}_{1i_0},$$

$$c_{i_0} = \frac{(\Gamma^{i_0}_{i_0j_0} - \Gamma^1_{i_0j_0})(T_{hi_0ij} + T_{hi_0ji} + T_{hii_0j} + T_{hiji_0} + T_{hji_0i} + T_{hjii_0})}{\Gamma^1_{ij} + \Gamma^1_{ji} - \Gamma^h_{ij} - \Gamma^h_{ji}} + S_{i_011i_0} - T_{1111},$$

$$\lambda^{i_0}_1 = \frac{-a_{i_0} - \sqrt{a_{i_0}^2 - 4c_{i_0}}}{2}, \quad \lambda^{i_0}_2 = \frac{-a_{i_0} + \sqrt{a_{i_0}^2 - 4c_{i_0}}}{2}.$$

The components $\{\Gamma_{ii}^h\}$ *of affine connection* ∇ *satisfy the following relations*

$$\begin{split} 2\Gamma_{ij}^{1} + \Gamma_{ji}^{1} &= 0, \quad 2\Gamma_{i_0j_0}^{\alpha} + \Gamma_{j_0i_0}^{\alpha} &= 0, \quad 2\Gamma_{ij}^{h} + \Gamma_{ji}^{h} &= 0, \quad 2\Gamma_{io}^{i_0} + \Gamma_{iio}^{i_0} &= 0, \\ \Gamma_{ij}^{1} + \Gamma_{ji}^{1} - \Gamma_{ij}^{k} - \Gamma_{ji}^{k} &= 0 \text{ for } k = i_0, \alpha, \quad 2\Gamma_{ik}^{k} + \Gamma_{ki}^{h} - 2\Gamma_{i1}^{1} - \Gamma_{1i}^{1} &= 0 \text{ for } k = i_0, h, \\ (\Gamma_{ij_0}^{1} + \Gamma_{ji_0}^{1} - \Gamma_{ij_0}^{h} - \Gamma_{joi}^{h}) S_{1i_011} &= 0, \quad (\Gamma_{ij_0}^{1} + \Gamma_{joi}^{1} - \Gamma_{ij_0}^{h} - \Gamma_{joi}^{h}) (2\Gamma_{io1}^{1} + \Gamma_{1i_0}^{1}) &= 0, \\ (\Gamma_{ij_0}^{1} + \Gamma_{ij_0}^{1} - \Gamma_{ij_0}^{h}) (S_{hij1} + S_{hji1}) + (\Gamma_{ij}^{1} + \Gamma_{ji}^{1} - \Gamma_{ij}^{h} - \Gamma_{ij_0}^{h}) T_{\alpha 111} &= 0, \\ (\Gamma_{ioj_0}^{\alpha} - \Gamma_{ioj_0}^{1}) (S_{hij1} + S_{hijol_0}) + (\Gamma_{ioj_0}^{1} + \Gamma_{ji_0}^{1} - \Gamma_{ij}^{h} - \Gamma_{joi_0}^{h}) T_{\alpha 111} &= 0, \\ (\Gamma_{ioj_0}^{\alpha} - \Gamma_{ioj_0}^{1}) (S_{kki1} + S_{kik1} - S_{1i11}) + (\Gamma_{ik}^{1} + \Gamma_{ki}^{1} - \Gamma_{ik}^{k} - \Gamma_{ki}^{k}) T_{\alpha 111} &= 0 \text{ for } k = i_0, h, \\ (\Gamma_{ij}^{1} + \Gamma_{ji}^{1} - \Gamma_{ij}^{h} - \Gamma_{ji}^{h}) T_{1ioj_0k} + (\Gamma_{ik}^{1} + \Gamma_{ki}^{1} - \Gamma_{ik}^{h} - \Gamma_{ki}^{h}) T_{1ioj_0j} + \\ (\Gamma_{ij}^{1} + \Gamma_{ji}^{1} - \Gamma_{ij}^{h} - \Gamma_{ij}^{h}) T_{1ioj_0k} + (\Gamma_{ik}^{1} + \Gamma_{ki}^{1} - \Gamma_{ik}^{h} - \Gamma_{ki}^{h}) T_{1ioj_0j} + \\ (\Gamma_{ij}^{1} + \Gamma_{ji}^{1} - \Gamma_{ij}^{h} - \Gamma_{ij}^{h}) (T_{\alpha ioj_0\alpha} - S_{11io_0}) + (\Gamma_{i\alpha}^{1} + \Gamma_{ia}^{1} - \Gamma_{i\alpha}^{h} - \Gamma_{\alpha i}^{h}) T_{\alpha ioj_0j} + \\ (\Gamma_{ij}^{1} + \Gamma_{ij}^{1} - \Gamma_{ij}^{h}) (T_{\alpha ioj_0\alpha} - S_{11io_0}) + (\Gamma_{i\alpha}^{1} + \Gamma_{ai}^{1} - \Gamma_{i\alpha}^{h} - \Gamma_{ai}^{h}) T_{\alpha ioj_0j} + \\ (\Gamma_{ij}^{1} + \Gamma_{ij}^{1} - \Gamma_{ij}^{h}) (T_{nioj_0} + T_{1ioj_0} + T_{1joi_0} + T_{1ioj_0} + T_{1ij_0j_0}) + \\ (\Gamma_{ij}^{1} - \Gamma_{ij}^{1} - \Gamma_{ij}^{h}) (T_{nioj_0} + T_{nio_0}^{1} + T_{nio_0}^{1} + T_{nij_0} + T_{nij_0} + T_{nij_0}) + \\ (\Gamma_{ij}^{1} + \Gamma_{ij}^{1} - \Gamma_{ij}^{h}) (T_{nioj_0}^{1} + T_{nio_0}^{1} + T_{nio_0}^{1} + T_{nij_0}^{1} + T_{nij_0}^{1} + T_{nij_0}^{1}) + \\ (\Gamma_{ij}^{1} + \Gamma_{ij}^{1} - \Gamma_{ij}^{h}) (T_{nioj_0}^{1} + T_{nio_0}^{1} + T_{nii_0}^{1} + T_{nij_0}^{1} + T_{nij_0}^{1} + T_{nij_0}^{1}) + \\ (\Gamma_{ij}^{1} + \Gamma_{ij}^{1} - \Gamma_{ij}^{1}) (T_{nioj_0}^{1} + T$$

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